Slide Lines in Lagrangian hydrocodes Combined Force and Velocity Boundary Condition

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Slide-line in Lagrangian hydro



Plan

- Context
 - 2D Lagrangian Staggered Hydrodynamics
 - Motivation for slide-line
 - Design principles
- Slide line method
 - Contact force
 - Velocity correction (slide-line vertex)
 - Improvements
- Numerical tests
 - Sanity checks
 - More advanced tests (Sedov, Rayleigh-Taylor)
- Beyond 2D and two-materials
- Conclusions and perspective

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Context

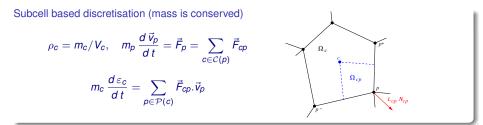
2D Staggered Lagrangian Hydrodynamics

2D Hydrodynamics in Lagrangian formulation

Euler hydrodynamics equations in an Inertial Confinement Fusion physics. Lagrangian formulation to follow multiple materials : large aspect ratio, shock waves, expansion, high deformation.

2D Compatible Staggered Lagrangian scheme

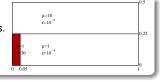
Thermodynamical variables at cell centers (ρ , ε), kinematics variables at vertexes (X, U). Conservative scheme (mass, momentum, total energy) on polygonal grid. Artificial viscosity to deal with shocks.

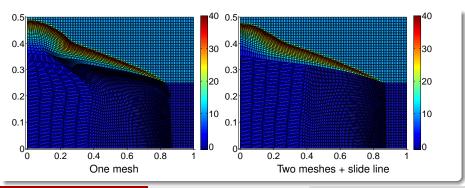




Context Motivation for slide line

Shear flows, vorticity are problematic for any Lagrangian methods. Example taken from [Caramana (JCP, 2009)]





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Context

Design principles and methods

Design principles

- Across a slide line normal velocity/acceleration is preserved, pressure is constant, tangential velocity is free.
- Conservation is ensured.
- No friction, no impact, no void opening. (Our choice !)
- Sliding treatment should not generate any inter-penetration or void opening. Robustness.
- Independence by respect to the meshes.
- Ideal situations must be perfectly solved.

Methods

Many different methods : [Wilkins (Academic Press, 1964)] : Classical approach with master/slave sides. [Benson (JCP, 1992)] : Review paper, followed by many contributors ! [Caramana (JCP, 2009)] : Approach compatible with our hydro scheme (based on forces, internal velocity BCs) — our starting point



Slide Line method

Contact force

Momentum equation for standard point p, no sliding :

$$m_{
ho}\,rac{d\,ec{v}_{
ho}}{d\,t}=ec{F}_{
ho}\,,$$

can be split for each separate half-point

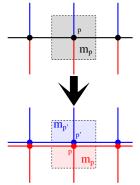
$$m_{
ho} \, rac{d \, ec{v}_{
ho}}{d \, t} = ec{F}_{
ho} + ec{g}_{
ho} \,, \ \ m_{
ho'} \, rac{d \, ec{v}_{
ho'}}{d \, t} = ec{F}_{
ho'} + ec{g}_{
ho'} \,.$$

- To be computed : contact forces g
 _p, g
 _{p'}, representing opposite pressure force.
- Suppose *p* and *p'* still coincide \Rightarrow same size, opposite orientation : $\vec{g}_{p'} = -\vec{g}_{p}$.
- Sum equations, same acceleration ⇒

$$ec{g}_{
ho} = -ec{g}_{
ho'} = rac{m_{
ho} \, ec{F}_{
ho'} - m_{
ho'} \, ec{F}_{
ho}}{m_{
ho} + m_{
ho'}}$$

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Slide Line method

Contact force

● Normal of each edge N_{p±1/2} used for construction of nodal normals N_p = (N_{p-1/2} + N_{p+1/2})/2.

• Normalization : $a_p = \|\vec{N}_p\|, \hat{c}_p = \vec{N}_p/a_p.$

- All integrals (present in masses and forces) must be scaled to lengths of opposite edges.
- Lower node : quantities from upper side must be scaled N_{p'-1/2} by factor a_p/a_{p'},

$$\vec{g}_{\rho} = \frac{m_{\rho} \frac{a_{\rho}}{a_{\rho'}} \vec{F}_{\rho'} - \frac{a_{\rho}}{a_{\rho'}} m_{\rho'} \vec{F}_{\rho}}{m_{\rho} + \frac{a_{\rho}}{a_{\rho'}} m_{\rho'}} = \frac{a_{\rho} \left(m_{\rho} \vec{F}_{\rho'} - m_{\rho'} \vec{F}_{\rho} \right)}{a_{\rho'} m_{\rho} + a_{\rho} m_{\rho'}}$$

- Upper node : scaling by $a_{p'}/a_p$, $\vec{g}_{p'} = -\frac{\frac{a_{p'}}{a_p} m_p \vec{F}_{p'} - m_{p'} \frac{a_{p'}}{a_p} \vec{F}_p}{\frac{a_{p'}}{a_p} m_p + m_{p'}} = \frac{a_{p'} \left(m_{p'} \vec{F}_p - m_p \vec{F}_{p'}\right)}{a_{p'} m_p + a_p m_{p'}}$
- Both equations same up to prime sign.

Design : Simple situations are perfectly solved independently of the meshes

 $\begin{array}{c|c} & a_{p}, \\ & a_{p}, \\ & & a_{p}, \\ & & \hat{c}_{p}, \\ & & \hat{c}_{p+1/2}, \\ & & \hat{a}_{p}, \\ & & \hat{a}_{p}, \end{array}$

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 $\hat{\mathbf{g}}_{n} \cdot \hat{\mathbf{c}}_{n}$

 \hat{c}_{p}

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Slide Line method

Contact force

- Contact forces only introduce interaction between sides due to pressures in each side ⇒ direction in the slide line normal.
- Projection of the contact forces to the normal direction,

$$\begin{split} m_p \frac{d \, \vec{v}_p}{d \, t} &= \vec{F}_p + (\vec{g}_p \cdot \hat{c}_p) \, \hat{c}_p, \\ m_{p'} \frac{d \, \vec{v}_{p'}}{d \, t} &= \vec{F}_{p'} + (\vec{g}_{p'} \cdot \hat{c}_{p'}) \, \hat{c}_{p'} \end{split}$$

Final projected contact forces :

Again same up to prime symbol.

Design : Normal component of acceleration is the same on both sides



Slide line method

Preventing Inter-Penetration by Velocity Correction

- Contact forces represent pressure (and other) forces from opposite side, no mechanism to prevent inter-penetration.
- From now, we need to start distinguishing between master and slave side, slave point (*p*, defined in initialization) must follow its master point :
 - $\vec{v}_{\rho}^{n+1} = \vec{v}_{\rho}^{n+1,\dagger} + \left[\left(\vec{v}_{\rho'}^{n+1,\dagger} + \vec{v}_{\rho'}^{n} \right) \cdot \hat{c}_{\rho'} \right] \hat{c}_{\rho'} \left[\left(\vec{v}_{\rho}^{n+1,\dagger} + \vec{v}_{\rho}^{n} \right) \cdot \hat{c}_{\rho'} \right] \hat{c}_{\rho'} \,.$
- Here, corrected \vec{v}_p^{n+1} is used for nodal motion, $\vec{v}_p^{n+1,\dagger}$ includes contact forces.
- Fix removes excessive velocity in the direction of inter-penetration (slide line normal) from the final velocity.
- Prevents slave node to move in this direction more than master does.
- Some sort of "internal velocity boundary condition".

Design : Inter-penetration or void opening are reduced, but conservation is lost.



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Goodness criteria for the slide line method

Measure of Energy Discrepancy

Total energy is lost !

- Change of internal energy from energy equation : m_c (εⁿ⁺¹_c − εⁿ_c) = − ∑_{p∈P(c)} F^c_p · Δr_p.
- Change of kinetic energy : $\frac{1}{2} m_p \left((\vec{v}_p^{n+1})^2 (\vec{v}_p^n)^2 \right)$.
- Compatible discretization ⇒ both values must be identical. However, due to velocity correction, energy discrepancy occurs.

Measure of Energy Discrepancy

- Discrepancy in point *p* during one timestep : $\Delta W_{p}^{n \to n+1} = \left[\frac{1}{2} m_{p} \left((\vec{v}_{p}^{n+1})^{2} - (\vec{v}_{p}^{n})^{2} \right) \right] - \left[\sum_{c \in C(p)} \vec{F}_{p}^{c} \cdot \Delta \vec{r}_{p} \right].$
- Integration over space and time :

$$\Delta W^{n \to n+1} = \sum_{\forall p} \Delta W_p^{n \to n+1}$$
, $\Delta W = \sum_{n=1}^N \Delta W_{n \to n+1}$.

Can be also generalized for non-compatible discretizations.

Check : Measure the impact of the slide line treatment.

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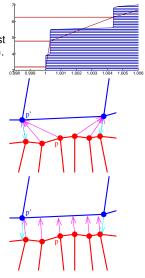


Improvement : interpolation of entity

- Up to now [Caramana (JCP, 2009)], interaction with closest point (in both contact force and velocity correction formulas).
- Produce disturbances of the slide line, especially in case of different aspect ratios, many points interact with only one
- All quantities on the opposite side must be interpolated :

$$\begin{array}{l} \left(\vec{g}_{\rho}\cdot\hat{c}_{\rho}\right)\,\hat{c}_{\rho}=-\frac{m_{\rho}\left(\vec{F}_{\rho'}\cdot\hat{c}_{\rho'}\right)+m_{\rho'}\left(\vec{F}_{\rho}\cdot\hat{c}_{\rho}\right)}{a_{\rho'}\,m_{\rho}+a_{\rho}\,m_{\rho'}}\,\,a_{\rho}\,\hat{c}_{\rho}\,,\\ \left(\vec{g}_{\rho'}\cdot\hat{c}_{\rho'}\right)\,\hat{c}_{\rho'}=-\frac{m_{\rho'}\left(\vec{F}_{\rho}\cdot\hat{c}_{\rho}\right)+m_{\rho}\left(\vec{F}_{\rho'}\cdot\hat{c}_{\rho'}\right)}{a_{\rho'}\,m_{\rho}+a_{\rho}\,m_{\rho'}}\,\,a_{\rho'}\,\hat{c}_{\rho'}\,,\\ \vec{v}_{\rho}^{n+1}=\vec{v}_{\rho}^{n+1,\dagger}+\left[\left(\vec{v}_{\rho'}^{n+1,\dagger}+\vec{v}_{\rho'}^{n}\right)\cdot\hat{c}_{\rho'}\right]\hat{c}_{\rho'}\\ -\left[\left(\vec{v}_{\rho}^{n+1,\dagger}+\vec{v}_{\rho}^{n}\right)\cdot\hat{c}_{\rho'}\right]\hat{c}_{\rho'}\,, \end{array}$$

i.e. normal \hat{c} , characteristic length a, mass m, force \vec{F} , and velocity \vec{v} .

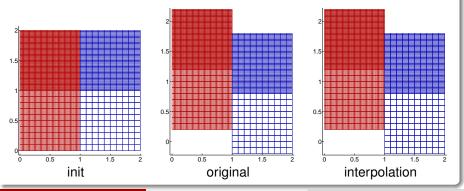




Sanity checks : Pure sliding

Vertically sliding blocks

- Non-uniform meshes, straight slide line, opposite velocity.
- Same pressures \Rightarrow contact forces must cancel exactly with pressure forces.
- Test of scaling.
- $\Delta W = 0$ for all methods, slide line perfect.



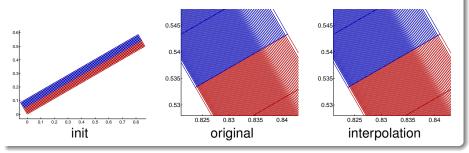
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Sanity checks : Salzmann-Like Piston

Oblique piston (from left to right)

- Sanity test : Salzmann piston with uniform mesh, straight slide line, rotated by π/6.
- Pressure gradient parallel with interface ⇒ contact forces cancel pressure forces.
- Test of scaling and projection.
- $\Delta W = 0$ for all methods, slide line perfect.

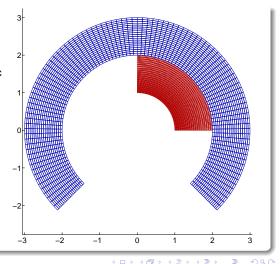


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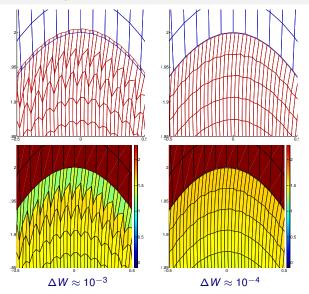
Numerical Examples Sliding Rings

- 2 meshes of 100 × 20 cells : static heavy outer master mesh (ρ = 10⁴) and moving light inner slave mesh (ρ = 1).
- Same pressure p = 1, T = 0.65 (original method fails soon after).
- Original method ⇒ serious distortions (driven by staircase problem), interpolation helps.



Slide-line in Lagrangian hydro





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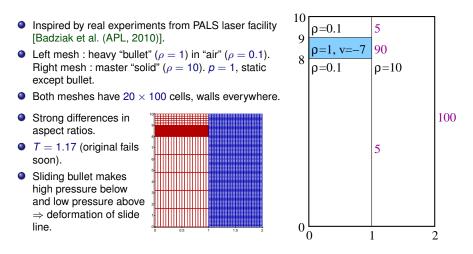
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Bullet in Channel



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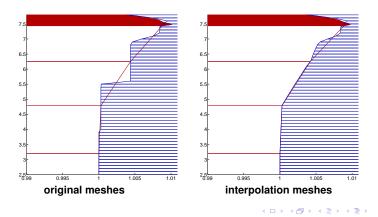
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Bullet in Channel

Early time T = 0.1

 $\bullet\,$ Slide line starts to be in motion \Rightarrow staircase problem for original method due to 1-to-1 interaction.



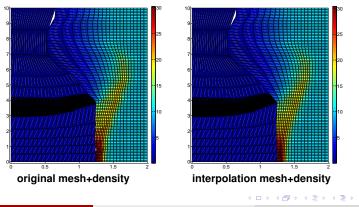
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Bullet in Channel

Later time T = 1.17

- Bullet and interface deformed.
- Original method soon to fail biggest inter-penetration and voids !



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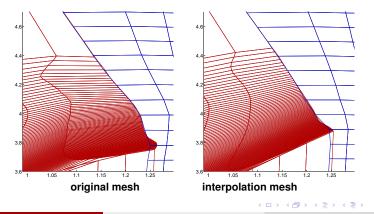
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Bullet in Channel

Later time T = 1.17

- Zoom on slide line at bullet.
- Bigger interpenetration and void for original method.



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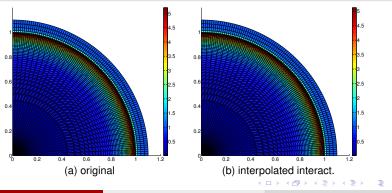
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Sedov problem with interface (slide-line)

Test symmetry violation due to slide-line machinery

- Sedov point explosion on polar mesh (unit density, zero pressure) at t = 1.
- Radius $r \in \langle 1/100, 1.1 \rangle$ and angle $\theta \in \langle 0, \pi/2 \rangle$
- Slide-line at r = 0.5, 20 cells in r direction, 100/31 outer/inner cells in angular direction. Meshes do not coincide



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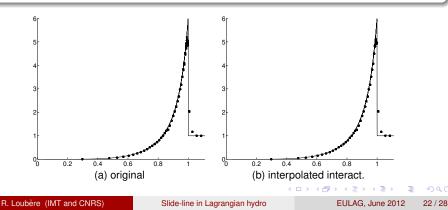
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Sedov problem with interface (slide-line)

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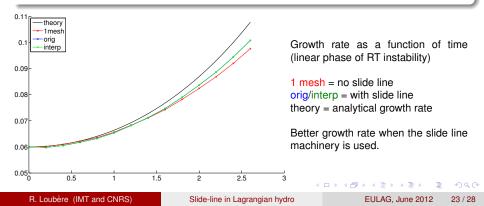




Numerical results

Rayleigh-Taylor Instability with Sliding

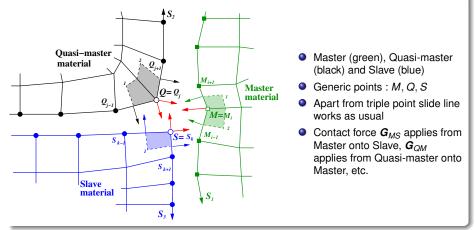
- Goal : analyse the influence of the slide line treatment on the growth rate of the Rayleigh-Taylor instability problem against the theoretical rate.
- Two simulations : 1 single 100 × 600 mesh (no sliding) and two 100 × 300 meshes (separated by a slide line)
- Upper side is the master side.





Beyond 2D and two-materials

Notation





Beyond 2D and two-materials Contact forces

Pair of contact forces

Momentum equations for triple point P in matrix form

$$\left(\begin{array}{c}m_{M}\\m_{Q}\\m_{S}\end{array}\right)\frac{d}{dt}\boldsymbol{U}_{P}=\left(\begin{array}{c}\boldsymbol{F}_{M}\\\boldsymbol{F}_{Q}\\\boldsymbol{F}_{S}\end{array}\right)+\left(\begin{array}{ccc}\boldsymbol{0}&\boldsymbol{G}_{MQ}&\boldsymbol{G}_{MS}\\\boldsymbol{G}_{QM}&\boldsymbol{0}&\boldsymbol{G}_{QS}\\\boldsymbol{G}_{SM}&\boldsymbol{G}_{SQ}&\boldsymbol{0}\end{array}\right).$$

Final form of pair of forces

$$\mathbf{G}_{MQ} + \mathbf{G}_{MS} = \frac{m_M(\mathbf{F}_Q + \mathbf{F}_S) - (m_Q + m_S)\mathbf{F}_M}{(m_Q + m_S) + m_M}, \\
 -\mathbf{G}_{MQ} + \mathbf{G}_{QS} = \frac{m_Q(\mathbf{F}_M + \mathbf{F}_S) - (m_M + m_S)\mathbf{F}_Q}{(m_M + m_S) + m_Q}, \\
 -\mathbf{G}_{MS} - \mathbf{G}_{QS} = \frac{m_S(\mathbf{F}_M + \mathbf{F}_Q) - (m_M + m_Q)\mathbf{F}_S}{(m_M + m_Q) + m_S}.$$

Need rescaling and velocity correction (Q+S constrained by M then S constrained by Q)

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Beyond 2D and two-materials Three dimensions

Planes of contact

- Point-to-point slide line treatment almost applies directly in 3D
- Compatible slide line treatment does not depend on dimensions
- Choice of contact normals is important
- Computer machinery is important

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Conclusion and Perspectives

Conclusions

- Slide lines from [Caramana (JCP, 2009)] implemented in 2D staggered compatible Lagrangian code. Apply to our scheme almost ideally.
- However some improvements : interpolation, numerical surface tension
- Numerical tests on gases show potentiality of the method
- Beyond 2-D, 2-mats \rightarrow (upcoming results)

Perspectives

- Material strength and impact of materials. Both are mandatory to validate such an approach.
- Full ALE with appropriate BCs
- Test in 3D



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