

Design and Configuration of MPAS for Deep-Atmosphere NWP and Geospace Applications

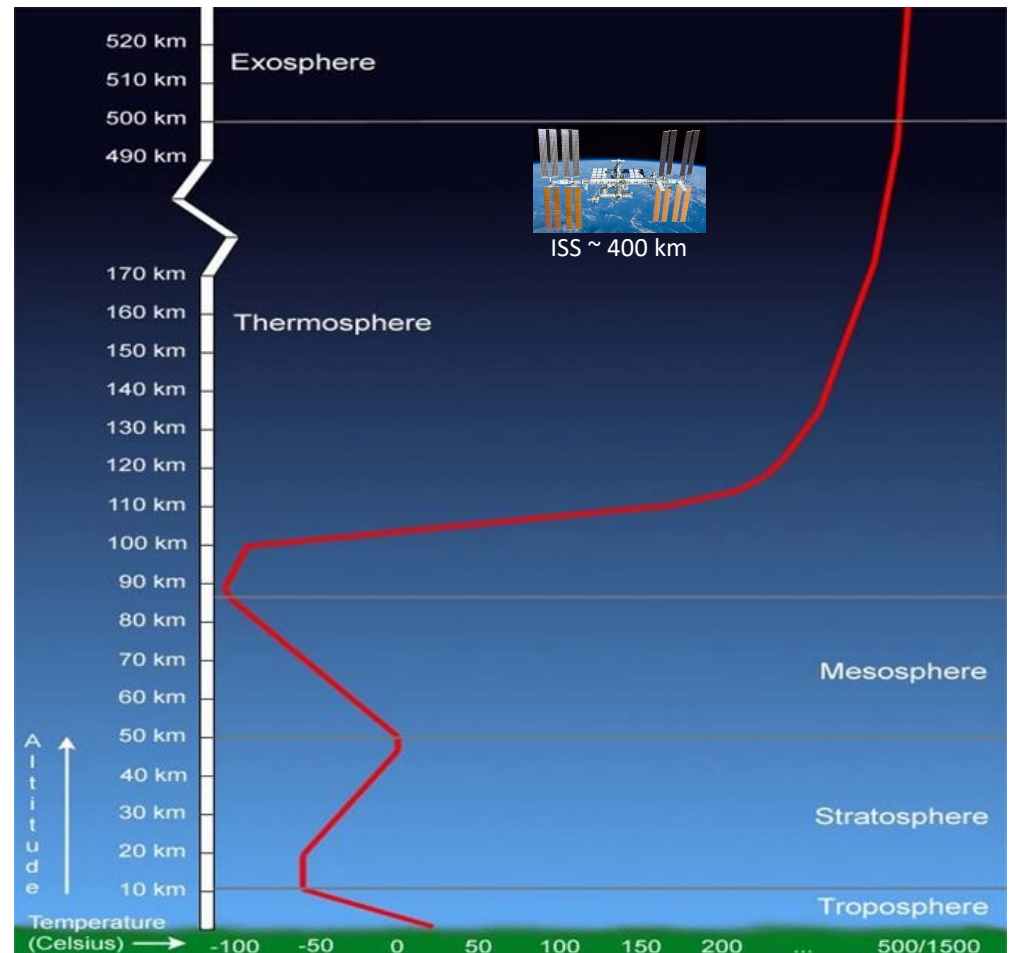
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Average Temperature Profile in the Earth's Atmosphere



10^{-8} 10^{-12}
 10^{-6} 10^{-9}
 10^{-4} 10^{-6}
 10^0 10^{-3}
 10^3 10^0 300
 10^{-5}

p (hPa) ρ (kg/m³) θ (K) v (m²/s)

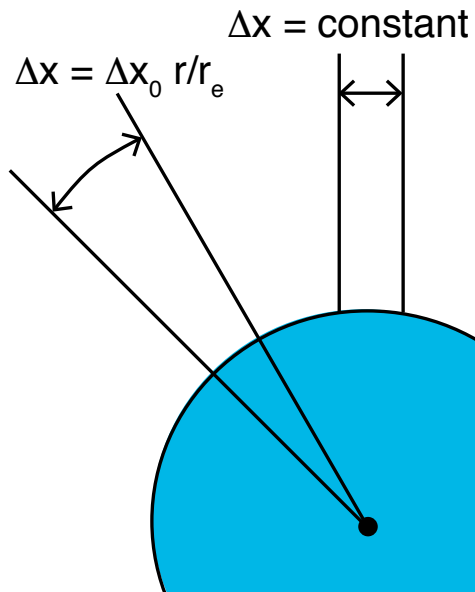
Deep Atmosphere (Full 3-D MPAS model):

- Use actual geocentric distance r instead of the Earth radius r_e in the governing equations and the grid mesh configuration
- Allow gravity to vary with height, $g(z) = g_0 \frac{r_e^2}{(r_e + z)^2}$
- Include Coriolis force terms involving vertical velocity components

Geospace, above ~ 150 km (2-D MPAS slab model):

- Include variable atmospheric composition and its coupling to the dynamics
- Modify the thermodynamic equation to account for variable composition
- Include (large) kinematic viscosity and thermal diffusivity terms
- Add appropriate physics and chemistry for the upper atmosphere (solar and Joule heating, ion drag, oxygen disassociation, etc.)

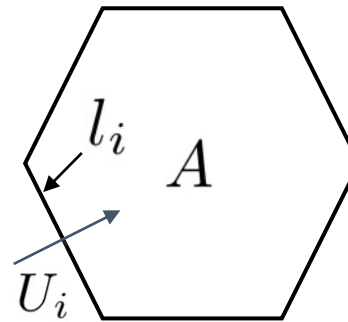
Deep Atmosphere Extensions for MPAS: Geometry



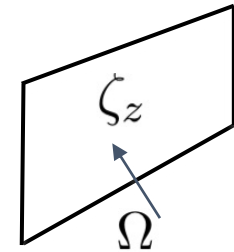
$$\frac{\partial \tilde{\rho}}{\partial t} = \nabla_{\zeta} \cdot \mathbf{V}_{\mathbf{H}} + \frac{\partial \Omega}{\partial \zeta}$$

$$\tilde{\rho} = \frac{\rho}{\zeta_z}, \quad \mathbf{V}_{\mathbf{H}} = (U, V) = (\tilde{\rho}u, \tilde{\rho}v), \quad \Omega = \tilde{\rho}\dot{\eta} \quad (\text{shallow atmosphere})$$

Hexagonal
grid cell:



horizontal



vertical

$$\frac{\partial \tilde{\rho}}{\partial t} = \frac{1}{A_0} \sum_{i=1}^n U_i l_{0i} + \frac{\Delta \Omega}{\Delta \zeta}$$

A_0, l_{0i} = Cell area and edge lengths at surface

Deep Atmosphere: $\tilde{\rho} = \frac{A}{A_0} \frac{\rho}{\zeta_z}, \quad U_i = \frac{A_0}{A} \frac{l_i}{l_{0i}} \tilde{\rho} u_i, \quad \Omega = \tilde{\rho} \dot{\eta}$

Shallow Atmosphere: $\tilde{\rho} = \frac{\rho}{\zeta_z}, \quad U_i = \tilde{\rho} u_i, \quad \Omega = \tilde{\rho} \dot{\eta} \quad (A = A_0, \quad l_i = l_{0i})$

Geospace Extensions for MPAS: Composition

Primary atmospheric constituents: $n = 3$ (N_2 , O_2 , O)

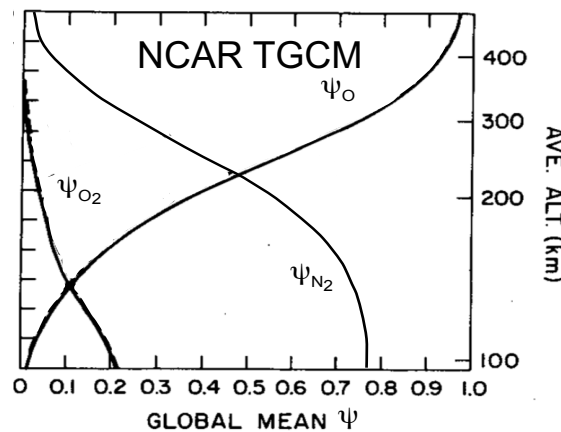
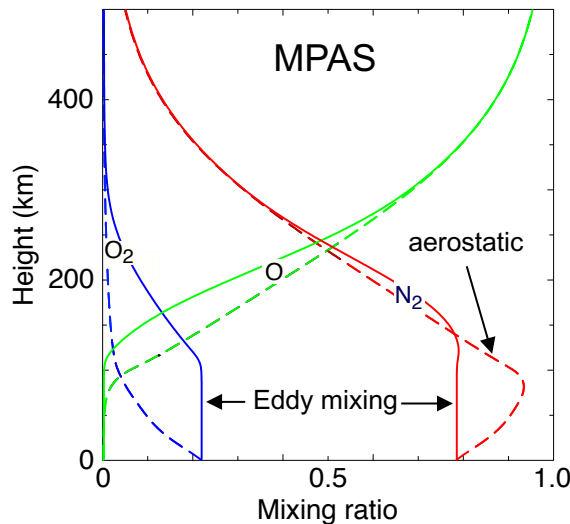
Aerostatic balance: constituents distributed according to their molecular weights:

$$\frac{\partial \ln p_i}{\partial z} = -\frac{gM_i}{k_b T}$$

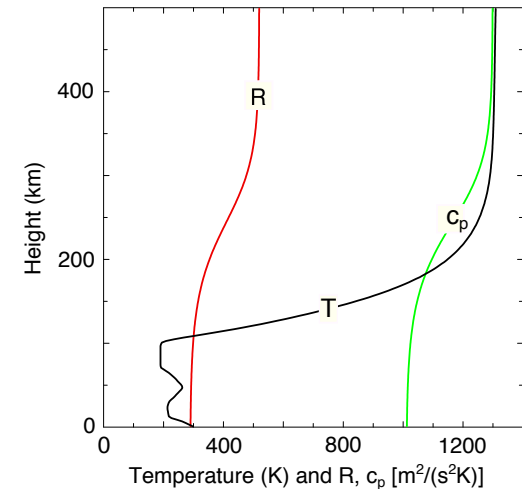
Thermodynamic Coefficients

$$\bar{R} = \frac{k_b}{\bar{M}} = k_b \sum_{i=1}^n \frac{\psi_i}{M_i}$$

$$\bar{c}_p = \sum_{i=1}^n c_{p_i} \psi_i$$



Dickinson et al. JAS, 1984



Composite thermodynamic properties vary only with the mixing ratios of the atmospheric constituents:

Geospace Extensions for MPAS: Thermodynamic Equation

Internal energy equation:
$$c_p \frac{dT}{dt} - RT \frac{d \ln p}{dt} = Q$$

From definition of potential temperature, $\theta = T p^{-\kappa}$ $\kappa = \frac{R}{c_p}$

$$\frac{d \ln \theta}{dt} = \frac{d \ln T}{dt} - \kappa \frac{d \ln p}{dt} - \ln p \frac{d \kappa}{dt}$$

$$\frac{d \ln \theta}{dt} = \frac{Q}{c_p T} - \ln p \frac{d \kappa}{dt} \quad (\neq 0 \text{ for } Q = 0)$$

Flux form thermodynamic equation:

$$\frac{\partial \Theta}{\partial t} + \nabla \cdot \mathbf{V} \theta = \frac{\rho Q}{c_p \pi} - \Theta \ln p \frac{d \kappa}{dt} \quad \Theta = \rho \theta, \quad \mathbf{V} = \rho \mathbf{v}$$

Small but potentially
significant correction !

$$\partial_t \mathbf{V}_H = -\mathbf{v}_H \nabla_\zeta \cdot \mathbf{V}_H - \partial_\zeta (\Omega \mathbf{v}_H) - \tilde{\rho}_d \nabla_\zeta K - p \left[\frac{\rho_d}{\rho_m} z_\zeta \nabla_\zeta \phi + z_H \frac{g z_\zeta}{R_d T_m} \right]$$

$$- \eta \mathbf{k} \times \mathbf{V}_H - \frac{\mathbf{v}_H W}{r} - e W \cos \alpha_r + (\nabla \tilde{\rho}_d \mathbf{v} \cdot \nabla) \mathbf{v}_H + F \mathbf{v}_H$$

$$\partial_t W = -(\nabla \cdot \mathbf{v} W)_\zeta - p \left[\frac{\rho_d}{\rho_m} \partial_\zeta \phi + \frac{g z_\zeta}{R_d T_m} \right]$$

$$+ \frac{uU + vV}{r} + e(U \cos \alpha_r - V \sin \alpha_r) + \nabla \tilde{\rho}_d \mathbf{v} \cdot \nabla W + F_W$$

$$\partial_t \Theta_m = -(\nabla \cdot \mathbf{V} \Theta_m)_\zeta - \Theta_m \phi \left[\frac{\partial \tilde{\rho}_d \kappa}{\partial t} + (\nabla \cdot \tilde{\rho}_d \kappa \mathbf{v})_\zeta \right] + \frac{1}{c_p \pi} \nabla \cdot (k_c \nabla T) + F_{\Theta_m}$$

$$\partial_t \tilde{\rho}_d = -(\nabla \cdot \mathbf{V})_\zeta$$

$$\partial_t Q_j = -(\nabla \cdot \mathbf{V} q_j)_\zeta + F_{Q_j},$$

$$\frac{\partial \Psi_i}{\partial t} = -\nabla \cdot \mathbf{V} \psi_i - \frac{\partial}{\partial \zeta} \left[\alpha^{-1} \mathbf{L} \psi \right] + \frac{\partial}{\partial \zeta} \left[\tilde{\rho}_d K_e(\zeta) \frac{\partial \psi_i}{\partial \zeta} \right] + \mathbf{S} - \mathbf{R}$$

$$\Theta_m = \tilde{\rho}_d \theta \left[1 + \frac{M_d}{M_v} q_v \right]$$

$$\kappa = \frac{R_d}{c_p}$$

$$p = p_0 \left(\frac{R_d \Theta_m}{p_0 z_\zeta} \right)^{\frac{1}{1-\kappa}}$$

$$\phi = \ln \frac{p}{p_o}$$

$$\eta = \mathbf{k} \cdot \nabla \times \mathbf{v}_H + f$$

$$e = 2\Omega_e \cos \psi$$

$$f = 2\Omega_e \sin \psi$$

Geospace Extensions for MPAS: Equations and Constituents

$$\begin{aligned}\partial_t \mathbf{V}_H &= -\mathbf{v}_H \nabla_\zeta \cdot \mathbf{V}_H - \partial_\zeta (\Omega \mathbf{v}_H) - \tilde{\rho}_d \nabla_\zeta K - p \left[\frac{\rho_d}{\rho_m} z_\zeta \nabla_\zeta \phi + z_H \frac{g z_\zeta}{R_d T_m} \right] \\ &\quad - \eta \mathbf{k} \times \mathbf{V}_H - \frac{\mathbf{v}_H W}{r} - e W \cos \alpha_r + (\nabla \tilde{\rho}_d \mathbf{v} \cdot \nabla) \mathbf{v}_H + F \mathbf{v}_H \\ \partial_t W &= -(\nabla \cdot \mathbf{v} W)_\zeta - p \left[\frac{\rho_d}{\rho_m} \partial_\zeta \phi + \frac{g z_\zeta}{R_d T_m} \right] \\ &\quad + \frac{uU + vV}{r} + e(U \cos \alpha_r - V \sin \alpha_r) + \nabla \tilde{\rho}_d \mathbf{v} \cdot \nabla w + F_W\end{aligned}$$

Coriolis terms
involving w

$$\partial_t \Theta_m = -(\nabla \cdot \mathbf{V} \Theta_m)_\zeta - \Theta_m \phi \left[\frac{\partial \tilde{\rho}_d \kappa}{\partial t} + (\nabla \cdot \tilde{\rho}_d \kappa \mathbf{v})_\zeta \right] + \frac{1}{c_p \pi} \nabla \cdot (k_c \nabla T) + F_{\Theta_m}$$

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Pressure gradients are
calculated using $\ln p$ instead of p

$$\Theta_m = \tilde{\rho}_d \theta \left[1 + \frac{M_d}{M_v} q_v \right]$$

$$\kappa = \frac{R_d}{c_p}$$

$$p = p_0 \left(\frac{R_d \Theta_m}{p_0 z_\zeta} \right)^{\frac{1}{1-\kappa}}$$

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 &\quad - \eta \mathbf{k} \times \mathbf{V}_H - \frac{\mathbf{v}_H W}{r} - e W \cos \alpha_r + (\nabla \tilde{\rho}_d \mathbf{v} \cdot \nabla) \mathbf{v}_H + F \mathbf{V}_H \\
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 \end{aligned}$$

Molecular
diffusivity terms

Thermal
Conductivity term

$$\begin{aligned}
 \partial_t \tilde{\rho}_d &= -(\nabla \cdot \mathbf{V})_\zeta \\
 \partial_t Q_j &= -(\nabla \cdot \mathbf{V} q_j)_\zeta + F_{Q_j},
 \end{aligned}$$

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 \end{aligned}$$

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Multiple equations are integrated for the atmospheric constituents, which determine R and c_p

$$\Theta_m = \tilde{\rho}_d \theta \left[1 + \frac{M_d}{M_v} q_v \right]$$

$$\kappa = \frac{R_d}{c_p}$$

$$p = p_0 \left(\frac{R_d \Theta_m}{p_0 z_\zeta} \right)^{\frac{1}{1-\kappa}}$$

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$$e = 2\Omega_e \cos \psi$$

$$f = 2\Omega_e \sin \psi$$

Only equations where variations
in atmospheric composition feed
back into the dynamics

Dynamical equations are solved using the same split-explicit
time integration as in MPAS for the shallow atmosphere

Geospace Extensions for MPAS: Mountain-wave test cases

2-D slab model

$$z_t = 500 \text{ km} \quad \Delta z = 2 \text{ km} \quad \bar{U} = 50 \text{ m/s}$$

$$h(x) = \frac{a^2 h_m}{(x - x_c)^2 + a^2} \quad a = 5\Delta x$$

$$\Delta x = 5, 20, 100 \text{ km} \quad h_m = 100 \text{ m}$$

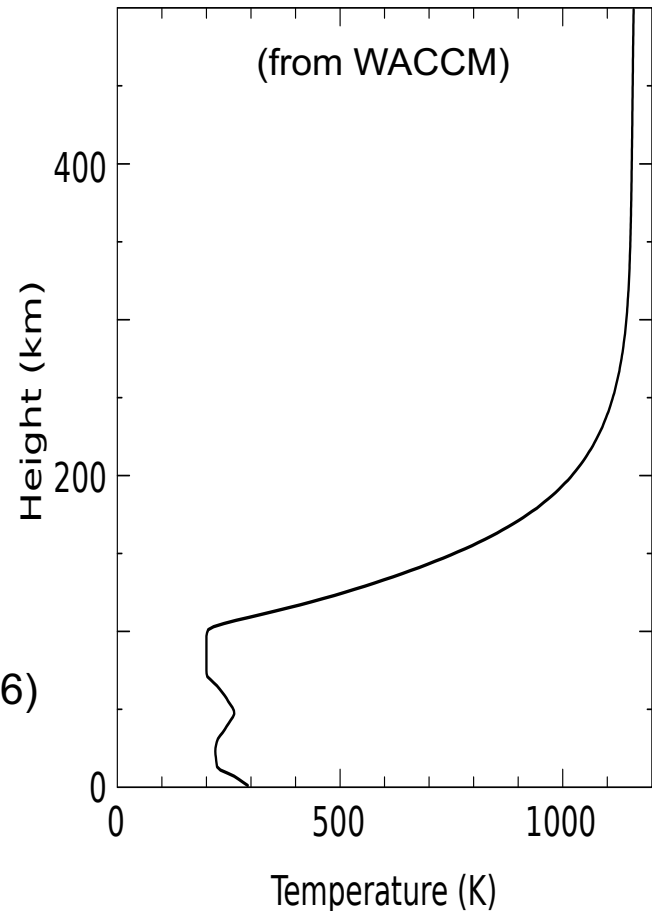
$$\Delta t = 4 \Delta x \quad \Delta \tau = \Delta t / 6$$

$$K_m = K_b + \ell^2 f(Ri) \frac{\partial u}{\partial z} \quad (\text{Hong et al. MWR 2006})$$

$$Ri = \frac{g}{\theta} \left[\frac{\partial \theta / \partial z}{(\partial u / \partial z)^2} \right]$$

Impulsive startup produces large amplitude acoustic modes in the upper domain. This startup noise is controlled by imposing Rayleigh damping on deviations from the initial state at early times.

Temperature Sounding



Mountain-wave test case *hydrostatic scales*

$$h_m = 100 \text{ m}$$

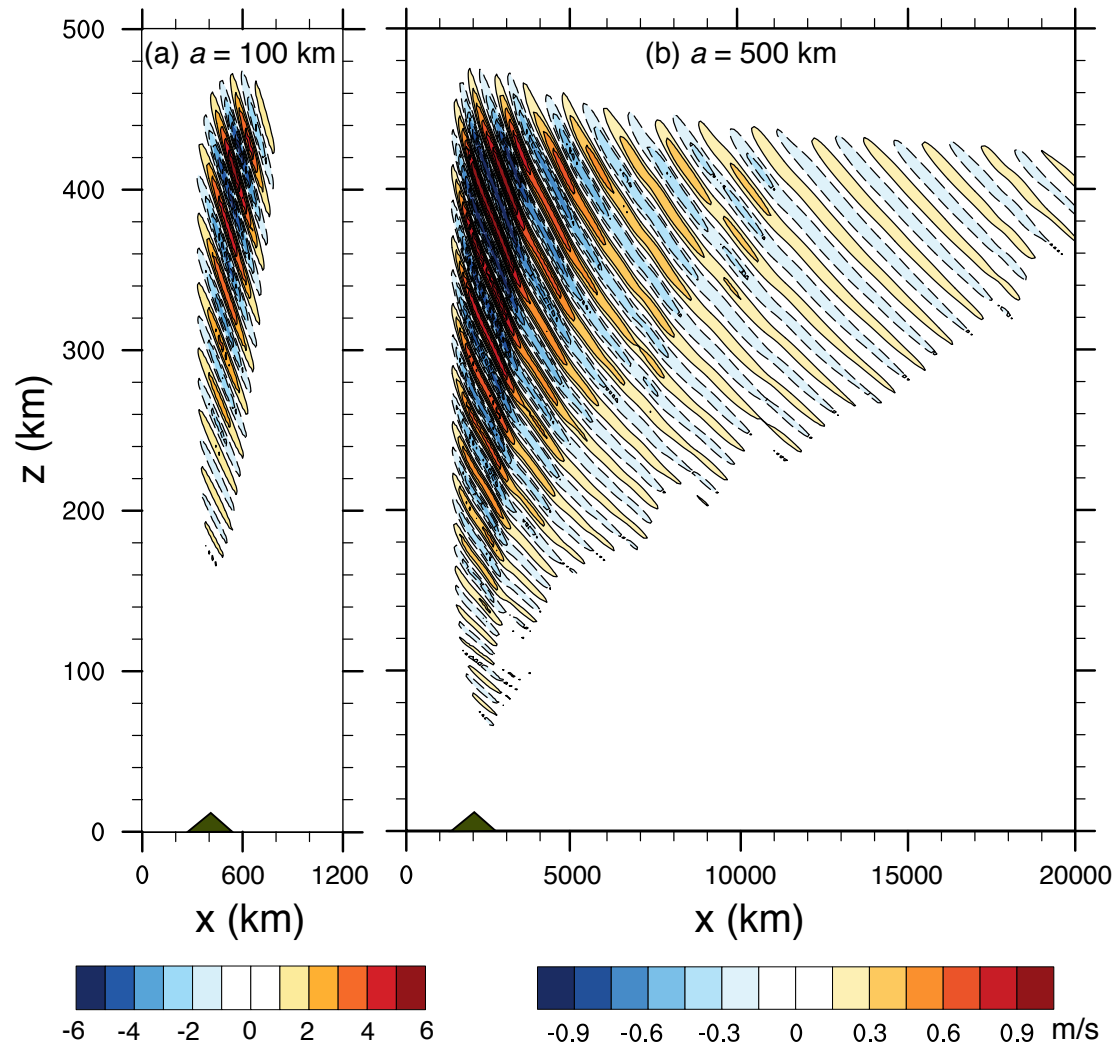
$$\Delta x = \text{(a) } 20 \text{ km,} \\ \text{(b) } 100 \text{ km}$$

$$K_b \Delta t / \Delta z^2 = \text{(a) } .03, \\ \text{(b) } .005$$

No molecular viscosity
or thermal conductivity

$$\ell^2 = k^2 \frac{N^2 - k^2 U^2}{k^2 U^2 - f^2}$$

Vertical velocity

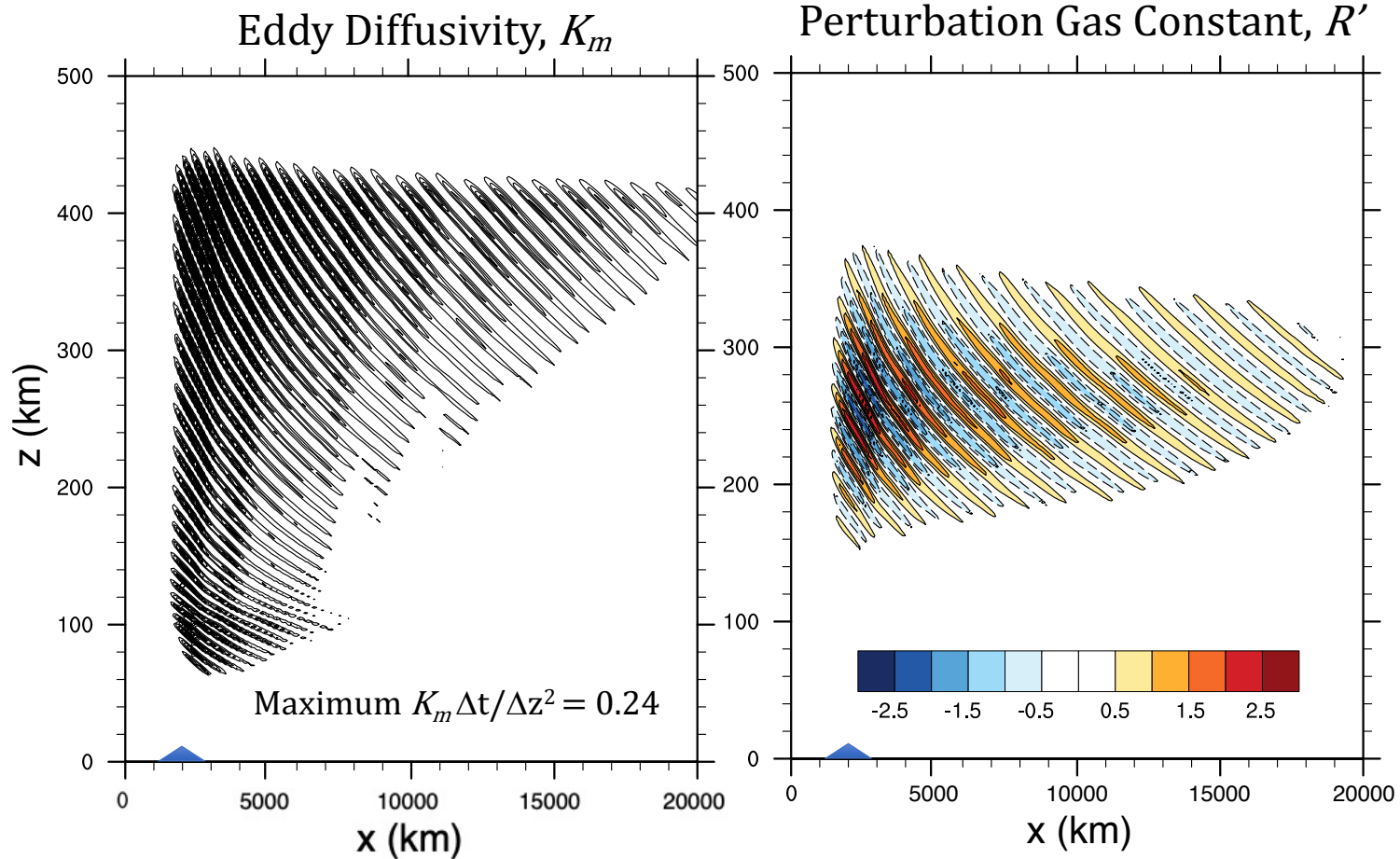


Mountain-wave test case *hydrostatic scales*

$$h_m = 100 \text{ m}$$

$$a = 500 \text{ km}$$

$$K_b \Delta t / \Delta z^2 = 0.005$$



No molecular viscosity or thermal conductivity

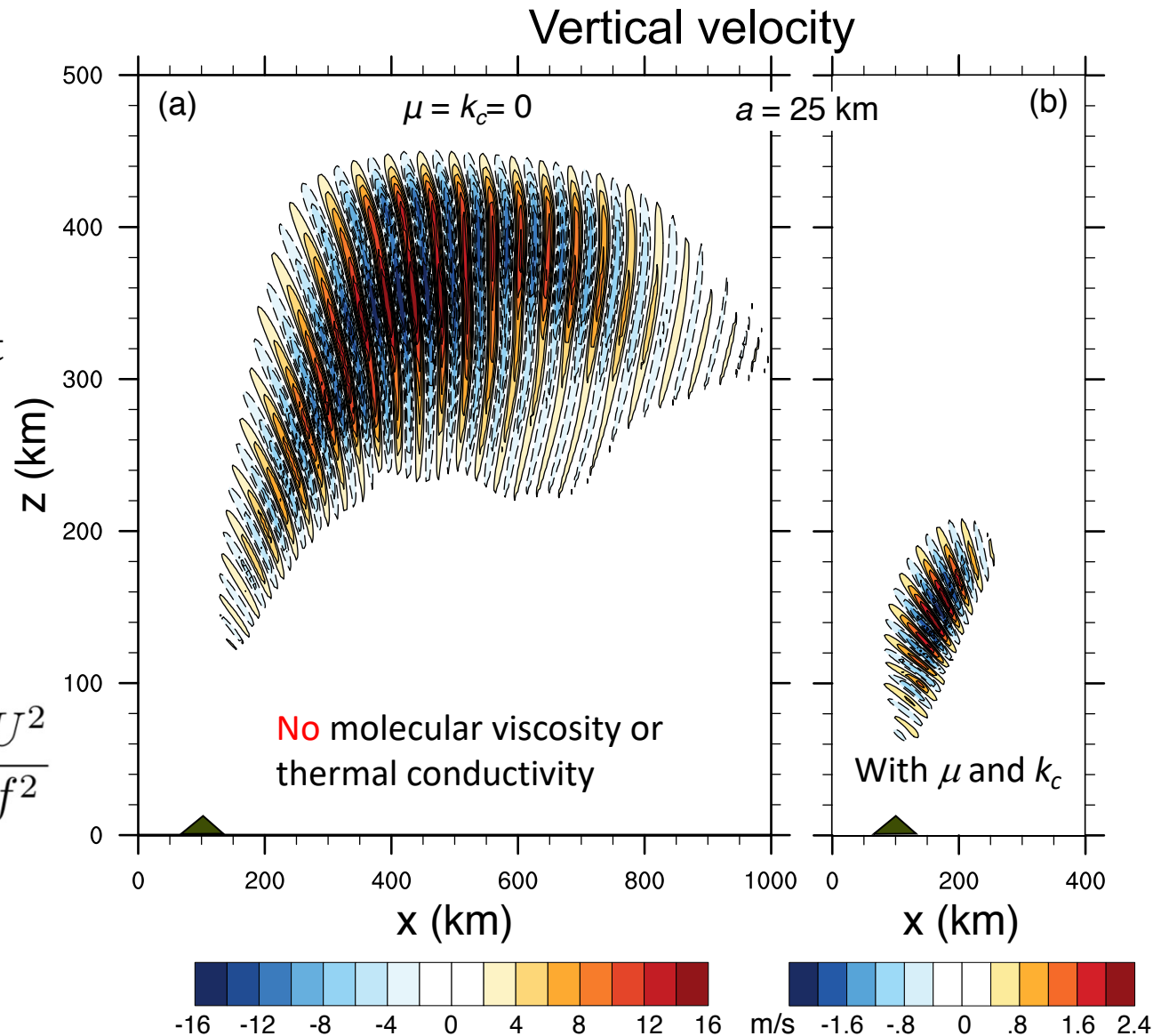
Mountain-wave test case *nonhydrostatic scales*

$$h_m = 100 \text{ m}$$

$$\Delta x = 5 \text{ km}$$

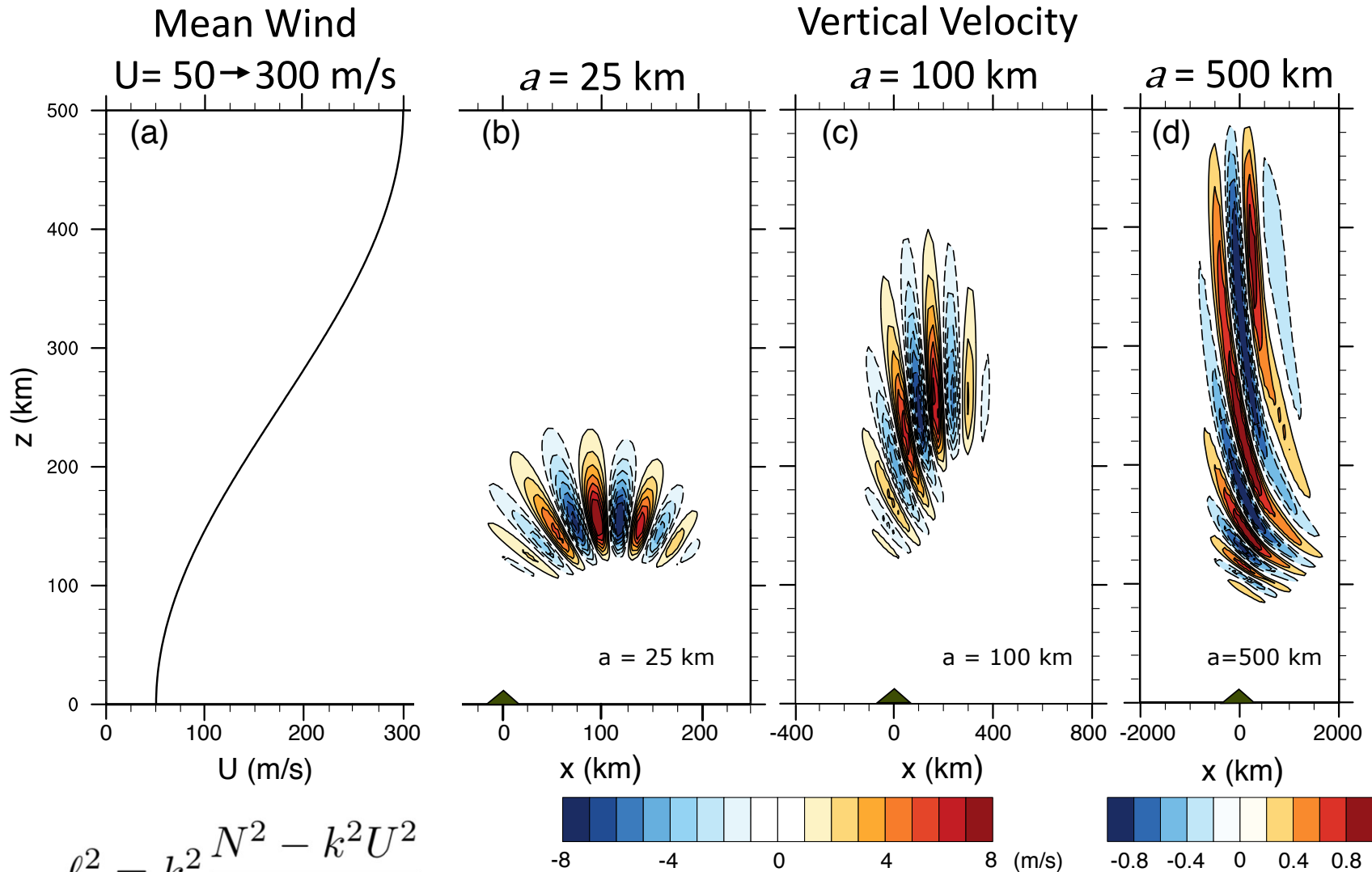
$$K_b = .027 \Delta z^2 / \Delta t$$

$$\ell^2 = k^2 \frac{N^2 - k^2 U^2}{k^2 U^2 - f^2}$$



Mountain-wave test case

Stronger mean wind



With molecular viscosity and thermal conductivity

- Height-based hybrid terrain-following coordinate seems well suited for deep atmosphere domains.
- Split-explicit finite-volume numerics for solving the nonhydrostatic dynamical equations appear to remain viable for the deep atmosphere.
- Equations are solved in double precision using full thermodynamic variables (not as perturbations from a reference state).
- Large molecular diffusivity and conductivity terms require special implicit treatment (currently using ADI).
- Large acoustic noise due to impulsive startup may require special filtering.
- Next steps: Implement in 3-D MPAS code, evaluate more realistic test cases, begin adding thermospheric physics.