

# **Implementation and early tests of a PDF parameterization in WRF**

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# Outline

1. Describe our subgrid-scale probability density function (PDF) parameterization
2. List disadvantages and advantages of PDF parameterizations
3. Show simulations of marine Sc clouds

# The cloud parameterization problem<sup>1</sup>

A parameterization of clouds and turbulence needs to supply subgrid-scale fluxes of heat, moisture, and momentum:

Heat:

$$\frac{\partial \bar{\theta}_l}{\partial t} = -\bar{w} \frac{\partial \bar{\theta}_l}{\partial z} - \frac{\partial}{\partial z} \overline{w' \theta'_l} + \bar{R} + \left. \frac{\partial \bar{\theta}_l}{\partial t} \right|_{ls} \quad (1)$$

Moisture:

$$\frac{\partial \bar{q}_t}{\partial t} = -\bar{w} \frac{\partial \bar{q}_t}{\partial z} - \frac{\partial}{\partial z} \overline{w' q'_t} + \left. \frac{\partial \bar{q}_t}{\partial t} \right|_{ls} \quad (2)$$

Momentum:

$$\frac{\partial \bar{u}}{\partial t} = -\bar{w} \frac{\partial \bar{u}}{\partial z} - f(v_g - \bar{v}) - \frac{\partial}{\partial z} \overline{u' w'} \quad (3)$$

$$\frac{\partial \bar{v}}{\partial t} = -\bar{w} \frac{\partial \bar{v}}{\partial z} + f(u_g - \bar{u}) - \frac{\partial}{\partial z} \overline{v' w'} \quad (4)$$

Red and Magenta = calculated by host model

Blue = calculated by parameterization

<sup>1</sup>According to Peter Stone of MIT.

# Our parameterization for clouds and turbulence: CLUBB

CLUBB denotes ``Cloud Layers Unified By Binormals.”

CLUBB is a unified, higher-order closure parameterization of subgrid variability. It predicts turbulence, shallow cumulus clouds, stratiform clouds, and deep convective clouds.

In our private copy of WRF, CLUBB replaces both PBL and deep convective schemes.

# CLUBB predicts a number of higher-order equations

In CLUBB, the set of prognosed moments includes:

$$\begin{aligned}
 \text{Means :} & \quad \frac{\partial \bar{u}}{\partial t} = \dots \quad \frac{\partial \bar{v}}{\partial t} = \dots \quad \frac{\partial \bar{q}_t}{\partial t} = \dots \quad \frac{\partial \bar{\theta}_l}{\partial t} = \dots \\
 \text{2nd - order :} & \quad \frac{\partial \overline{w'q'_t}}{\partial t} = \dots \quad \frac{\partial \overline{w'\theta'_l}}{\partial t} = \dots \quad \frac{\partial \overline{w'^2}}{\partial t} = \dots \\
 & \quad \frac{\partial \overline{q_t'^2}}{\partial t} = \dots \quad \frac{\partial \overline{\theta_l'^2}}{\partial t} = \dots \quad \frac{\partial \overline{q'_t\theta'_l}}{\partial t} = \dots \\
 \text{3rd - order :} & \quad \frac{\partial \overline{w'^3}}{\partial t} = \dots
 \end{aligned}$$

$w$  = vertical velocity       $q_t$  = total water specific humidity

$\theta_l$  = liquid water potential temperature

These moments predict the fluxes we need and some extra quantities as well.

# CLUBB closes many terms using an assumption about the shape of the subgrid probability density function (PDF)

CLUBB use the *Assumed* PDF Method. We *assume* a *functional form* of the PDFs, and determine a *particular instance* of this functional form for each grid box and time step. (The form we assume is a double Gaussian PDF.)

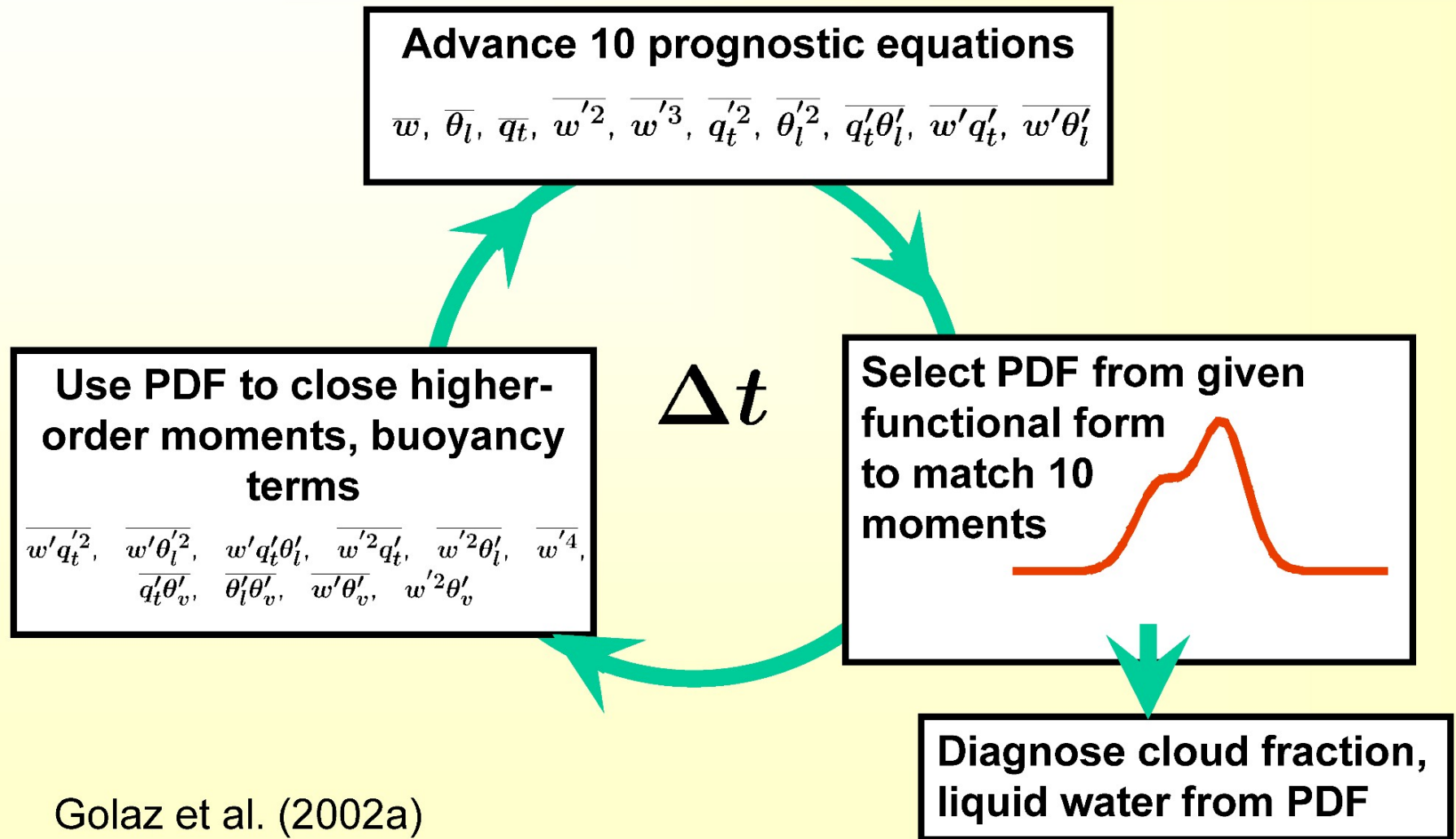
The PDF varies in space and evolves in time.

# Steps in the Assumed PDF Method

The Assumed PDF Method contains **3** main steps that must be carried out for each grid box and time step:

- (1) Prognose grid box means and various higher-order moments.
- (2) Use these moments to select a particular PDF instance from the assumed functional form.
- (3) Use the selected PDF to compute average of various higher-order terms that need to be closed, e.g. buoyancy flux, cloud fraction, etc.

## Schematic of the Assumed PDF method



Golaz et al. (2002a)



**What are some *disadvantages* of the assumed PDF method?**

## **Disadvantage 1 of the PDF approach: Complexity**

The code base for a PDF parameterization is larger than that for a mass-flux scheme.

The problem of complexity can be ameliorated to some extent by good software engineering.

## **Disadvantage 2 of the PDF approach: Computational cost**

The PDF method prognoses many higher-order moments. Therefore, it is expensive.

Optimization of the code can reduce the runtime of PDF parameterizations to some extent.

**What are some potential *advantages* of the assumed PDF method over other types of cloud parameterization, e.g. mass-flux schemes?**

## **Parameterization problem: Lack of memory**

Some convective parameterizations respond instantly to instability (e.g. CAPE). That is, the parameterizations are diagnostic, rather than prognostic.

In nature, however, deep convection evolves gradually from shallow convection. Lack of memory of whether shallow convection has occurred is part of the reason that deep convection sometimes onsets prematurely in parameterizations.

# The PDF approach: Including memory

All of CLUBB's higher-order moment equations contain a time-tendency term.

Therefore, CLUBB need not respond instantaneously to changes in CAPE. This allows the possibility of gradual onset of deep convection.

# **Parameterization problem: Inappropriate formulation in the terra incognita**

At grid spacings between 1 and 10 km, convection is partly resolved and partly subgrid scale. Therefore, these grid spacings are "terra incognita" for parameterizations.

Many parameterizations assume that cloudy updrafts occupy a small fraction of a grid column, but this assumption is no longer true at finer grid spacings.

(Wyngaard 2004; Gerard 2007)

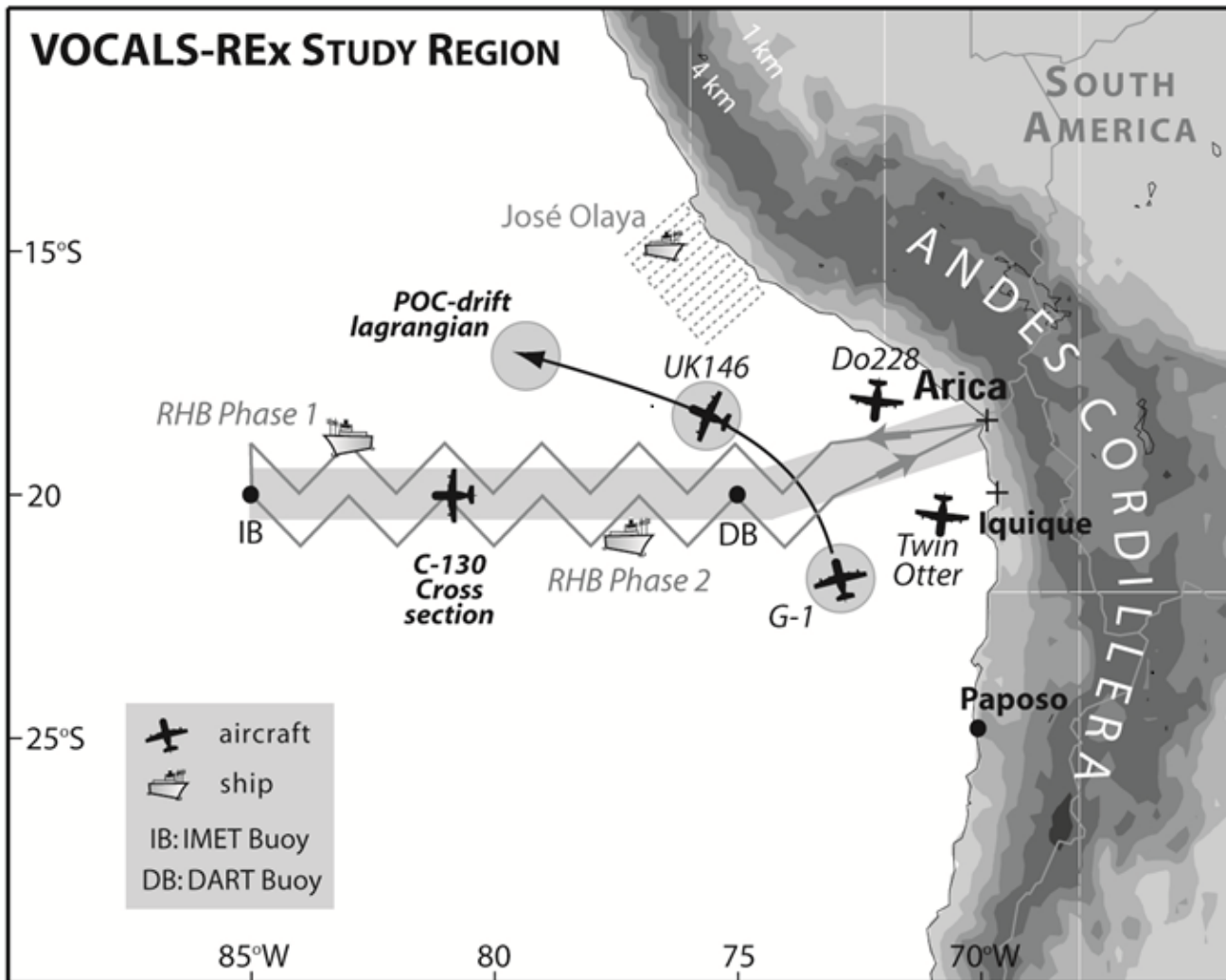
# The PDF approach: Formulation in the terra incognita

Because CLUBB does not employ a mass-flux scheme, it needs to make no assumption about the fraction of a grid box that is occupied by updrafts or environmental air.

From a PDF point of view, the key difference between small and large grid boxes is that in nature, smaller volumes encompass less variance of cloud and turbulence.



# A test case: The VOCALS field experiment studied marine Sc that occurred near deep convection over land



Courtesy  
Rob Wood

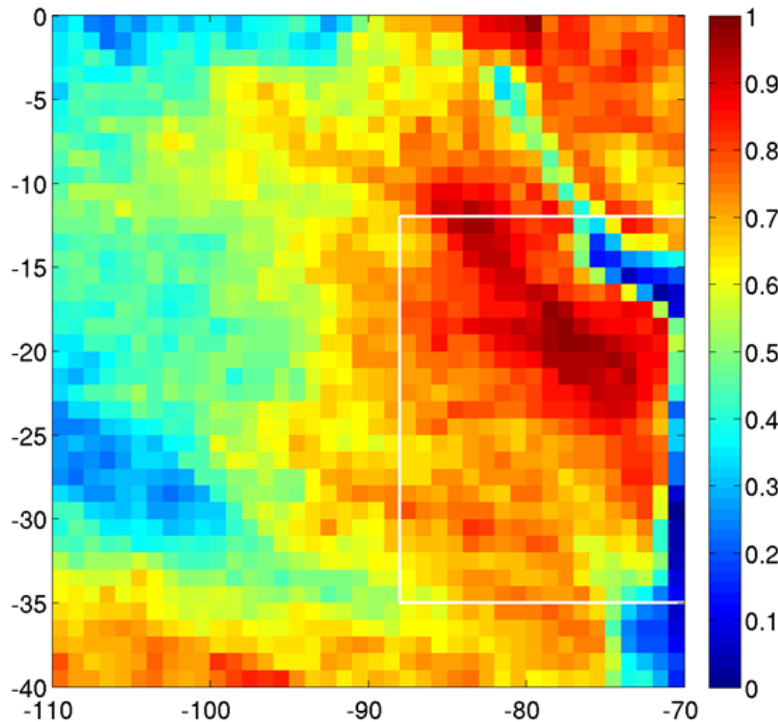
# **WRF-CLUBB is participating in Matt Wyant's model intercomparison of the VOCALS clouds**

- We simulate a month-long period during Oct/Nov 2008.
- In the plots that follow, we use coarse (50-km) grid spacing and a one-minute timestep.
- One deficiency of these simulations is that they don't (yet) feed cloud fraction from CLUBB to the microphysics.

# WRF-CLUBB produces Sc cloud, but places the cloud maximum too far north

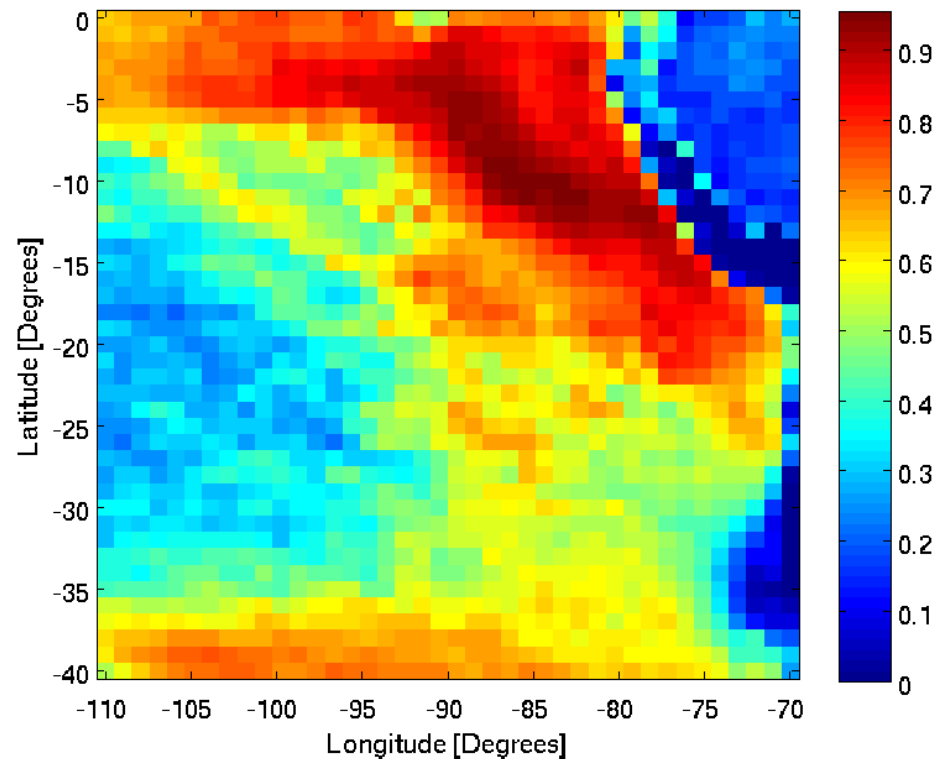
## MODIS satellite obs

Cloud cover



## WRF-CLUBB

Cloud cover [Dimensionless]



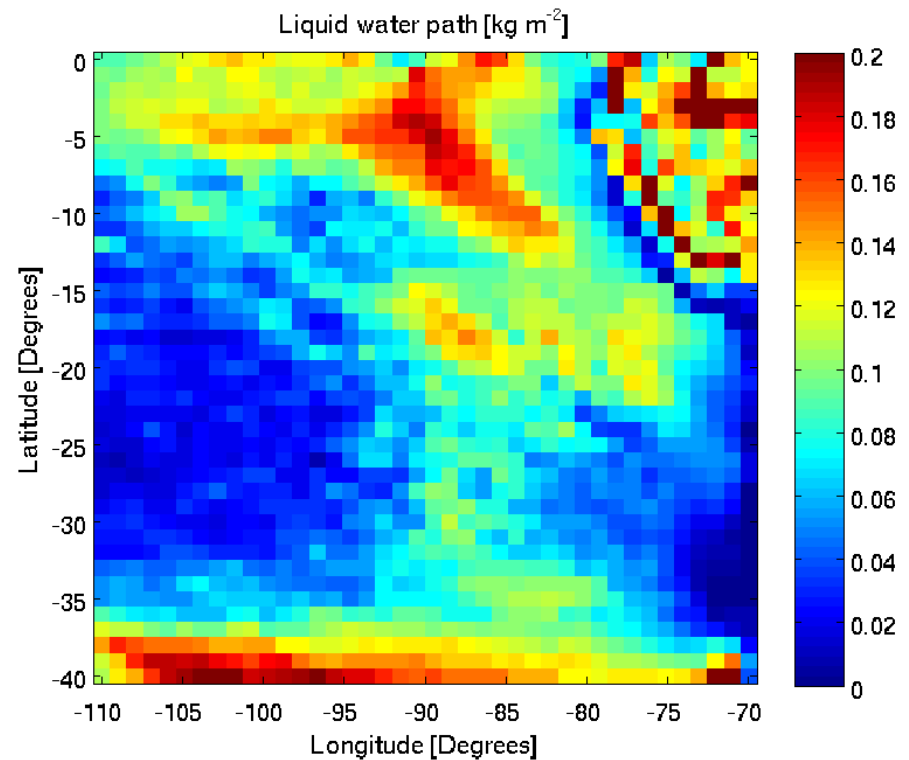
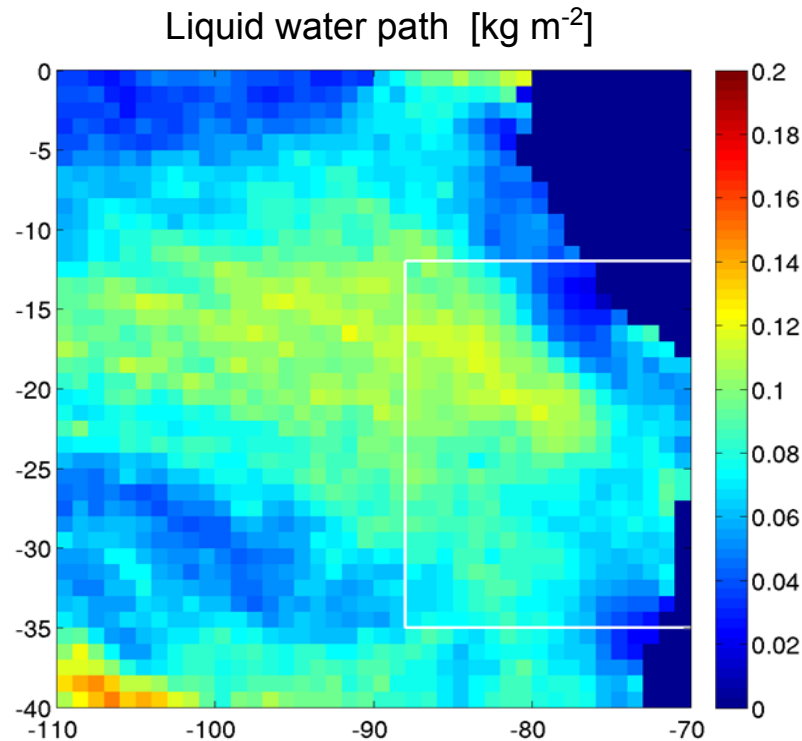
Courtesy Matt Wyant and Rob Wood

run035, 50km, GFS forcings, l\_mp\_clubb\_pdf=false

# WRF-CLUBB produces too much liquid to the north and south

**MODIS**

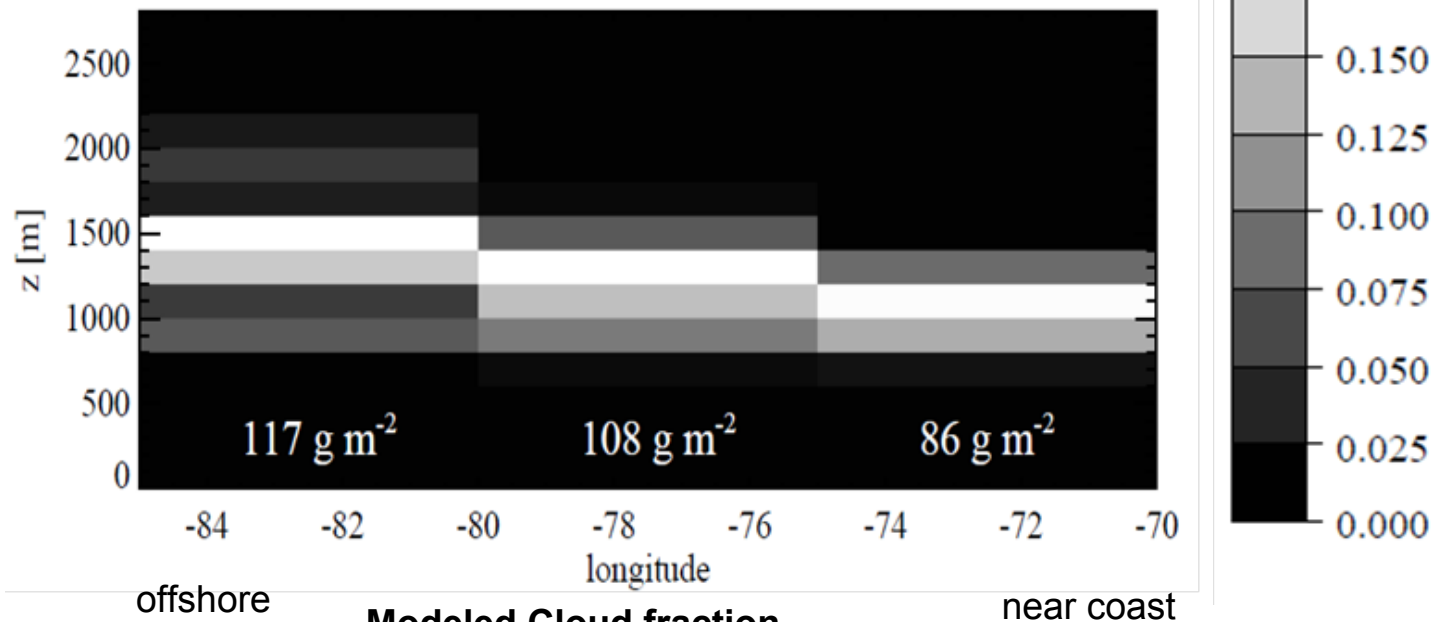
**WRF-CLUBB**



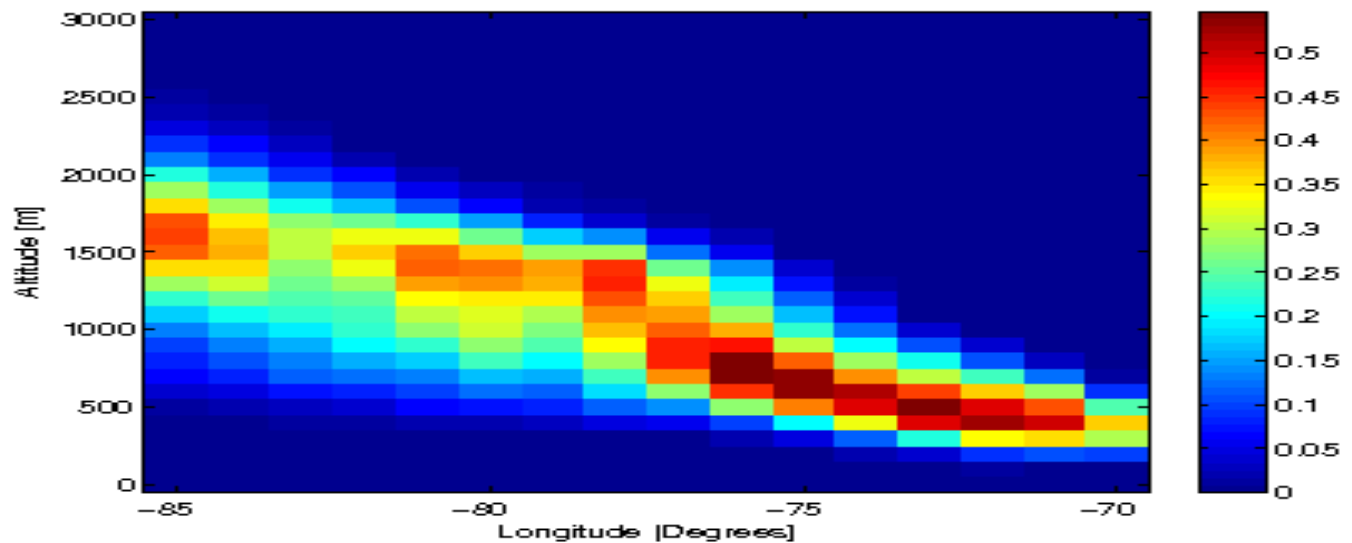
run035, 50km, GFS forcings, l\_mp\_clubb\_pdf=false

**WRF-CLUBB  
produces  
clouds that  
deepen  
offshore, but  
WRF-CLUBB's  
cloud-top  
height is too  
low near the  
coast**

**Observed cloud cross section along 20  
degrees south**

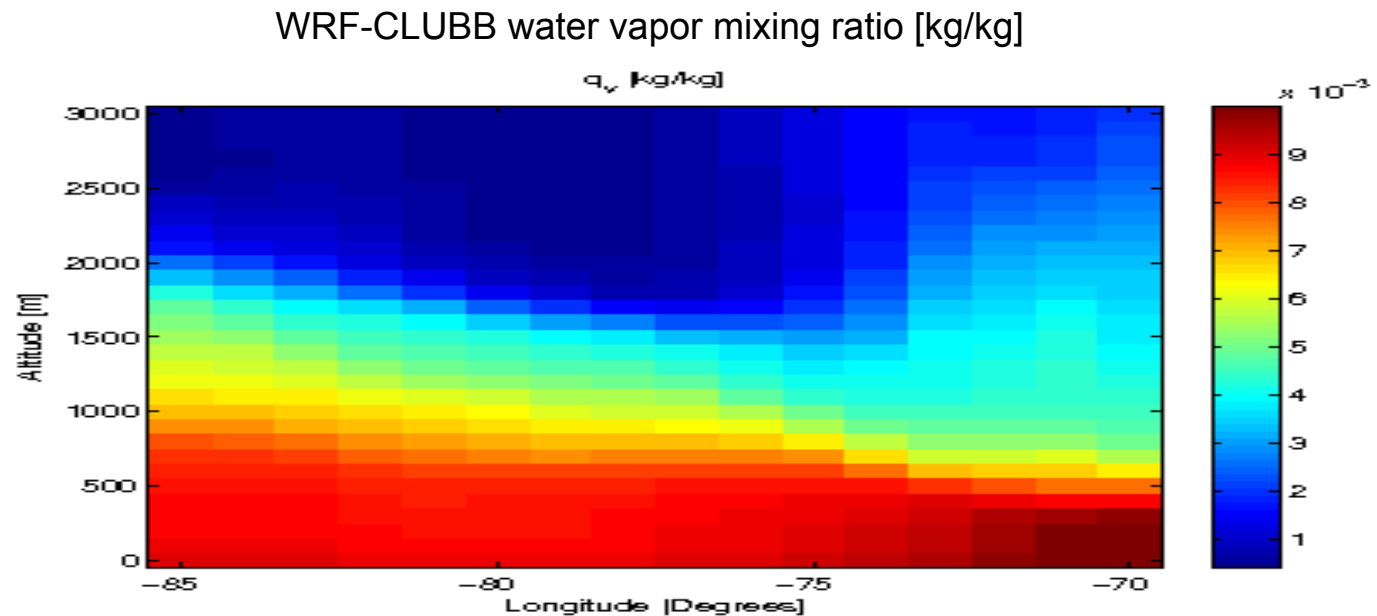
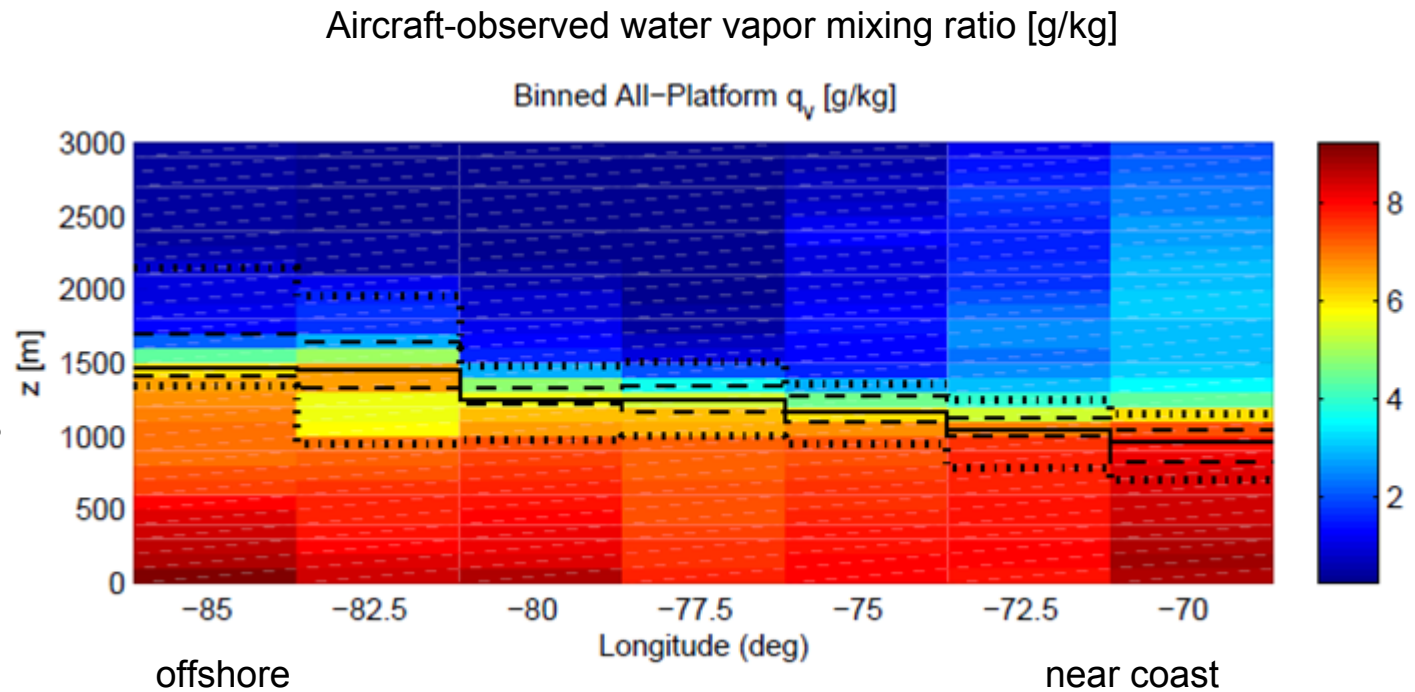


**Modeled Cloud fraction**



run035, 50km, GFS forcings, l\_mp\_clubb\_pdf=false

**WRF-CLUBB's  
cloud-top  
water vapor  
gradient isn't  
as sharp as  
observations**



run035, 50km, GFS forcings, l\_mp\_clubb\_pdf=false

# Conclusions

- We've implemented a PDF parameterization ("CLUBB") in WRF
- We've tested the parameterization for marine Sc. The simulations still need work, but nonetheless the results are encouraging.

Thanks for your time!



## **Parameterization problem: Inconsistent results as the horizontal grid spacing is refined**

Ideally, cloud parameterizations would produce similar results when used at different grid spacing. Insensitivity to resolution is needed especially for simulations with nested grids (e.g., Warner and Hsu 2000).

# **The PDF approach: Vary dissipation with grid box size**

Making the length scale a function of the grid spacing is a simple way of progressively reducing CLUBB's fluxes as resolution increases. As the grid spacing approaches the scale at which convection is resolved, CLUBB lets the resolved scales take over, as desired (Larson et al. 2012).